

Presentation/Publication Information: APS Meeting, March 11-16, 2001, Seattle, WA

NanoScience Colloquium, May 18-20

Acknowledgments:

This work was supported under NASA contract [RTOP: 519-40-12(NAS2-14303)].

Abstract:

Modeling Ballistic Current Flow in Carbon Nanotube Wires

M. P. Anantram

Experiments have shown carbon nanotubes (CNT) to be almost perfect conductors at small applied biases [1]. The features of the CNT band structure, large velocity of the crossing subbands and the small number of modes that an electron close to the band center / Fermi energy can scatter into, are the reasons for the near perfect small bias conductance. We show that the CNT band structure does not help at large applied biases - electrons injected into the non crossing subbands can either be Bragg reflected or undergo Zener-type tunneling. This limits the current carrying capacity of CNT [2]. We point out that the current carrying capacity of semiconductor quantum wires in the ballistic limit is different, owing to its band structure.

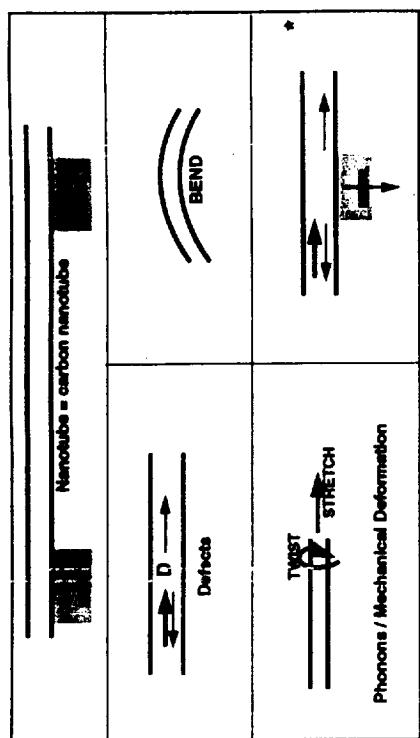
The second aspect addressed is the relationship of nanotube chirality in determining the physics of metal-nanotube coupling [3]. We show that a metallic-zigzag nanotube couples better than an armchair nanotube to a metal contact. This arises because in the case of armchair nanotubes, while the π band couples well, the π^* band does not couple well, to the metal. In the case of zigzag nanotube both crossing modes couple reasonably well to the metal. Many factors such as the role of curvature, strain and defects will play a role in determining the suitability of nanotubes as nanowires. From the limited view point of metal-nanotube coupling, we feel that metallic-zigzag nanotubes are preferable to armchair nanotubes.

[1] S. Frank, P. Poncharal, Z. L. Wang and W. A. de Heer, Conductance quantization in multi-walled carbon nanotube, *Science*, vol. 280, p. 1744 (1998); P. G. Collins, M. Hersam, M. Arnold, R. Martel, and Ph. Avouris, Current Saturation and Electrical Breakdown in Multiwalled Carbon Nanotubes *Phys. Rev. Lett.*, v. 86, p. 3128 (2001)

[2] M. P. Anantram, Current Carrying Capacity of Carbon Nanotubes *Phys. Rev. B*, vol. 62, p. R4837, (2000)

[3] M. P. Anantram, Which nanowire couples better to metals: Armchair or zigzag carbon nanotubes? *Appl. Phys. Lett.*, v. 78, p. 2055, (2001)

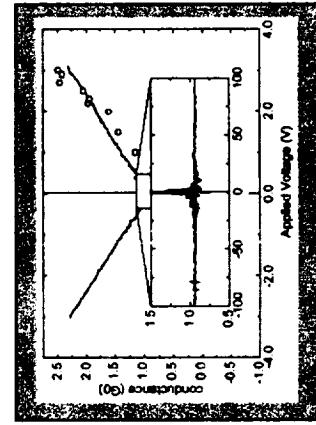
Topics Studied



* Bragg reflec.: Intrinsic mechanism, which exists even in an ideal situation

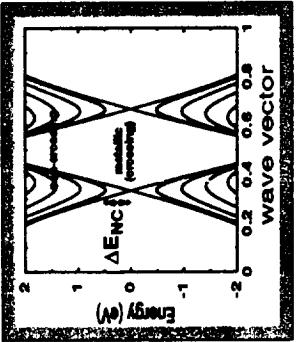
* This Poster

Frank et. al, Science 280 (1998)



- $V_{APPLIED} < 200mV$, $G \sim 2e^2/h$

- $V_{APPLIED} > 200mV$, slow increase



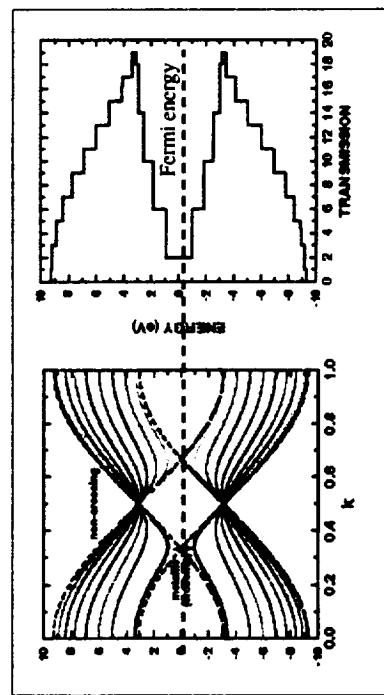
- $E \sim \pm 120meV$, non-crossing bands open
- At $E \sim 2eV$ electrons are injected into about 80 subbands

- Yet the conductance is only $\sim 3.75 e^2/h$

Modelling Ballistic Current Flow in Carbon Nanotube Wires

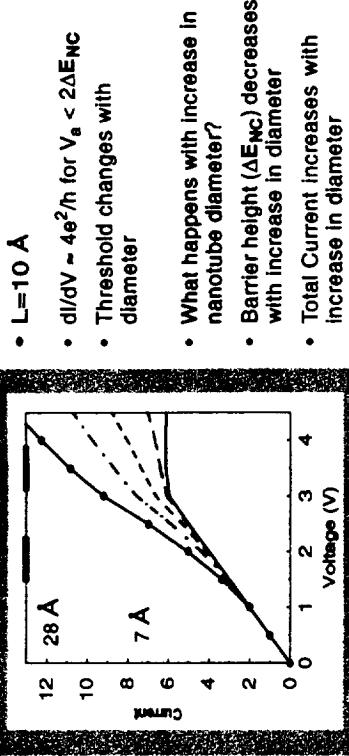
M. P. Anantram
NASA Ames Research Center
Moffett Field
CA 94035-1000

Current-carrying capacity of carbon nanotubes

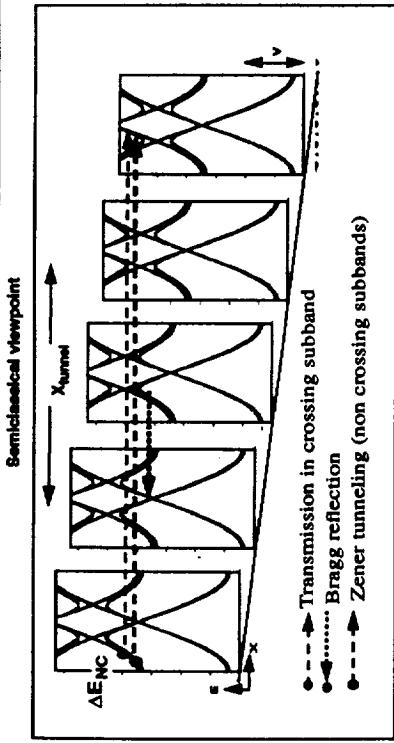


- Close to $E=0$, only two sub-bands, Conductance = $\frac{4e^2}{h}$ ($6k\Omega$)
- At higher energies, Conductance = $\frac{(20-30)e^2}{h}$ ($< 1k\Omega$)

Can subbands at the higher energies be accessed to drive large currents (small resistance) through these molecular wires?



The differential conductance is NOT comparable to the increase in the number of subbands.
For a (20,20) nanotube, there are 35 subbands at $E = \pm 3.5V$.



• The strength of the two processes are determined by:

- Tunnelling distance (X_{tunel}) \rightarrow Screening length
- Barrier height, $2\Delta E_{NC}$
- Scattering and Defects

$\Delta E_{NC} \propto 1/\text{Diameter}$. So, the importance of Zener tunneling increases with increase in nanotube diameter.

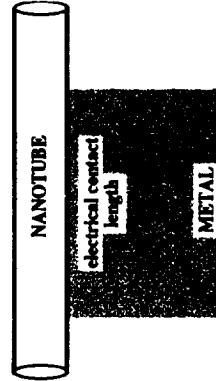
Coupling of carbon nanotubes to metallic contacts

Electronic properties of nanotubes are closely related to chirality:

- Metal versus Semiconductor
- Bandgap change with deformation / strain.

Questions:

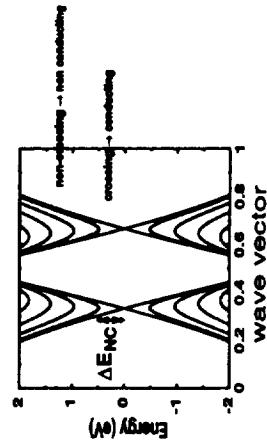
Is there a preferable nanotube chirality to maximize current flow?
Role of wave vector conservation?
Explain experimentally observed scaling of conductance with contact length



Parameters that influence current flow:

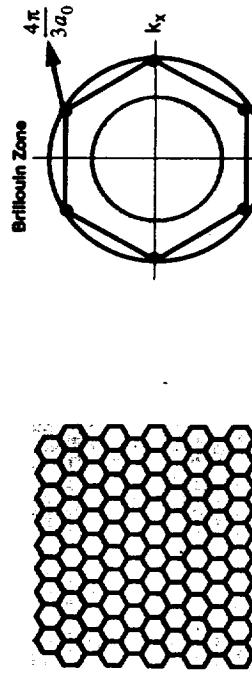
- Strength of coupling to metal
- Length of metal-nanotube contact
- Defects
- Metal Fermi wave vector

Bragg reflection severely limits the current carrying capacity
The crossing metallic-type bands conduct current.
Current carrying capacity of non-crossing subbands is limited.



Large diameter nanotubes: non-crossing bands will partially conduct due to Zener-type tunneling.
Conductances significantly larger (ten times) than $4e^2/h$ would be difficult.

GRAPHENE SHEET IN UNIFORM CONTACT WITH METAL



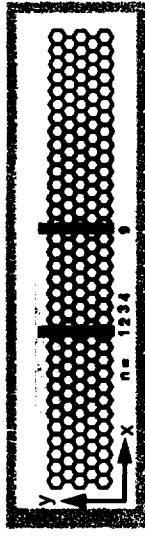
FIG

For good coupling: Metal $k_F \text{Fermi} > 4\pi/3a_0$ (1.7 Å)

$k_F \text{Fermi}$ Å ⁻¹	
Cs	0.65
Ag	1.20
Au	1.21
Hg	1.37
O	1.75
	1.7

Ashcroft & Mermin, Solid State Physics J. Tersoff, Appl. Phys. Lett. v. 74, p. 2122 (1999)

Scattering rate



$$\Psi = e^{ink_x L} \phi \quad n = \text{integer and } \phi \text{ is wave func. of atoms in a 1D unit cell}$$

Scattering rate from metal to nanotube (Born approx.):

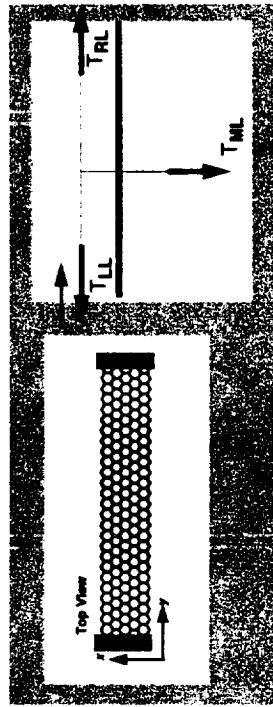
$$1/\tau \alpha | < \Psi_{nt} | V_{m-nt} | \Phi_m > |^2$$

$$\delta(k_x - k_x^m) \quad | < \phi | V_{m-nt} | \Phi_m > |^2$$

- K_x is conserved
- K_y conservation is relaxed due to finite width of contact area

How do we model the system?

- π electron tight binding model
- Metal is modelled as a free electron gas (k_F)

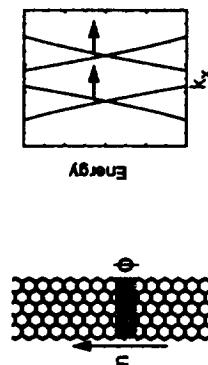


$$T_{RL} + T_{ML} + T_{LL} = 2.$$

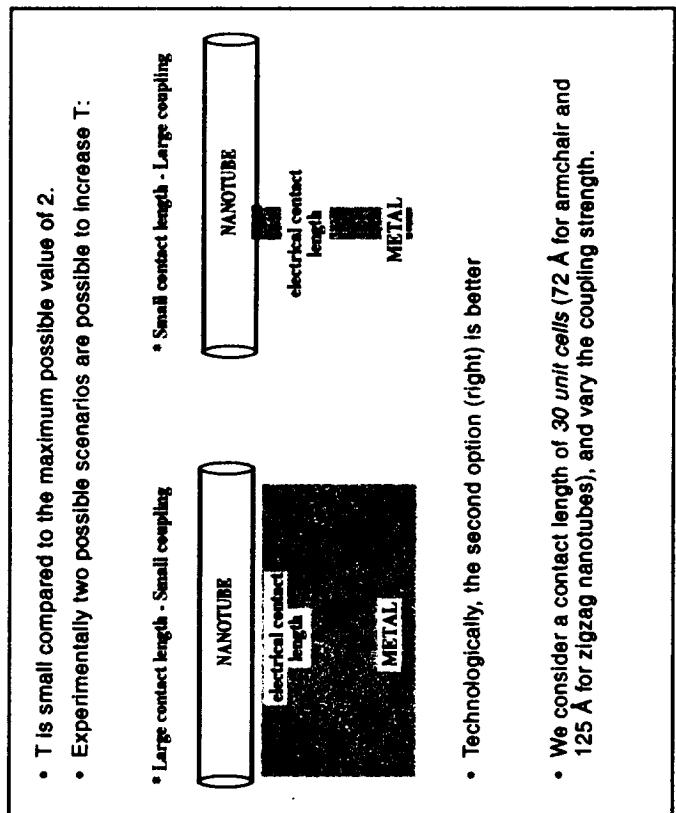
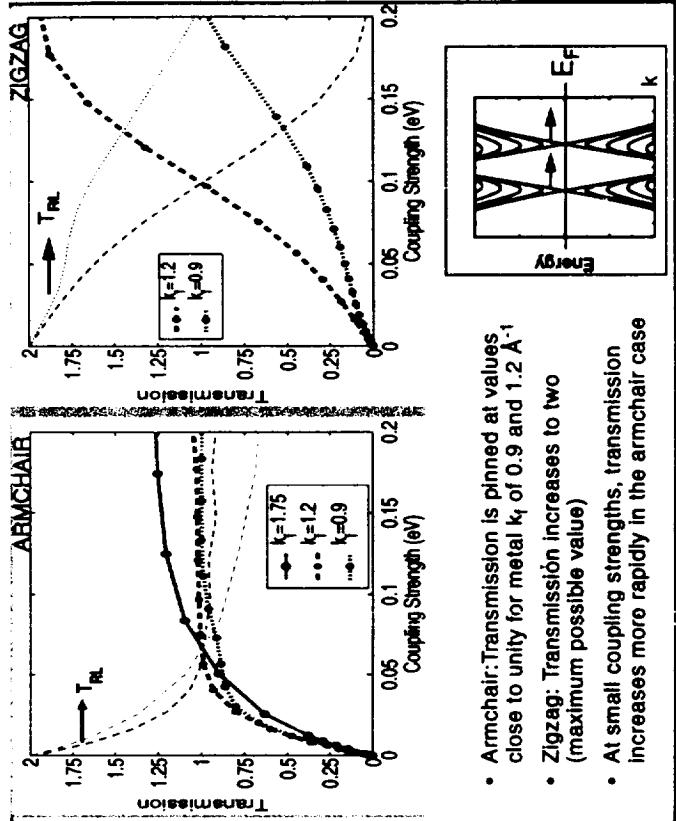
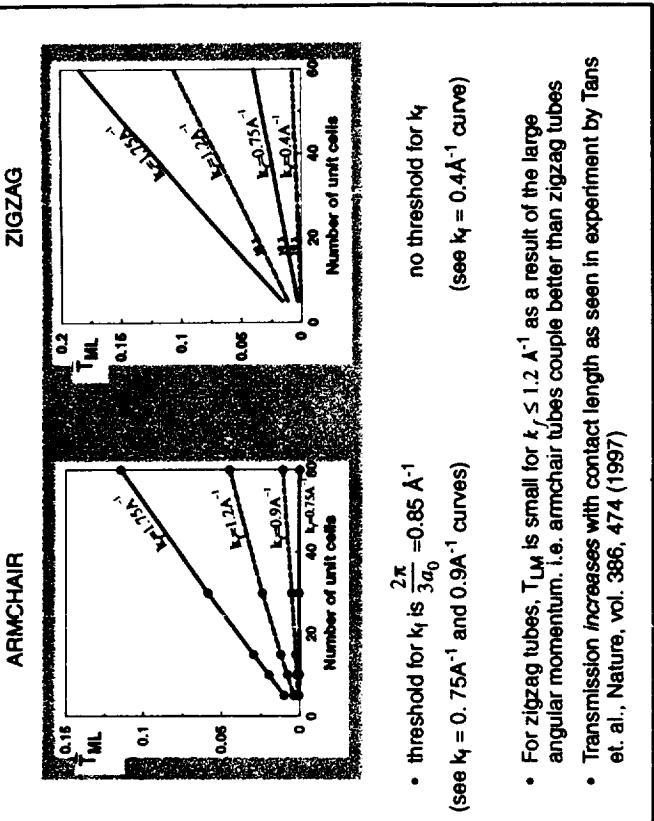
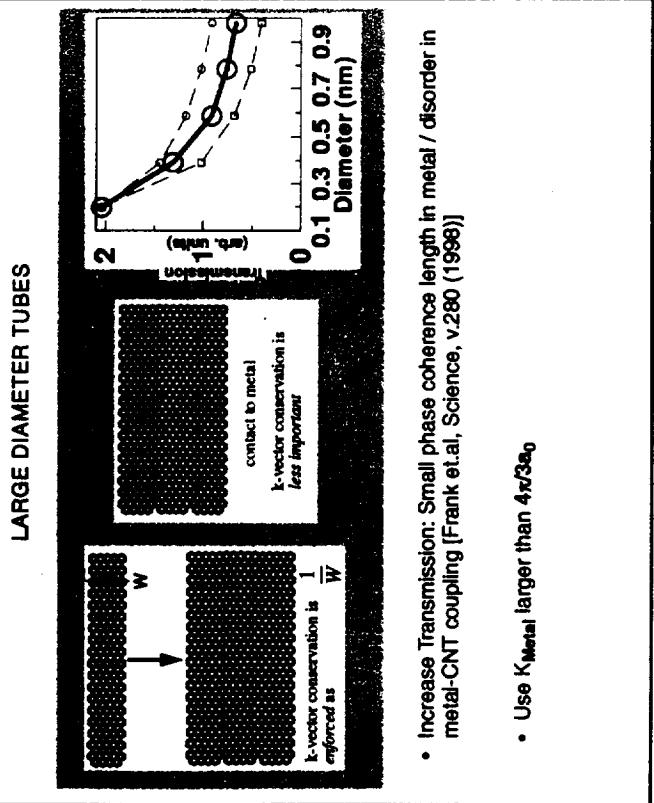
Below threshold $k_F \text{Fermi} - T$ does not scale with increase in contact length.

- Phys. Rev. B v.58, p. 4882 (1998) and v. 61, p. 14219 (2000)
- Compute self energy due to: (i) metal & (ii) semi-infinite CNT leads

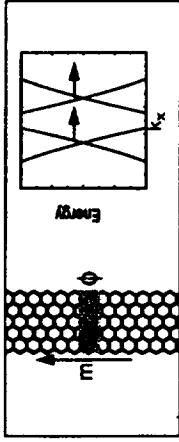
$$1/\tau \alpha \quad \delta(k_x - k_x^m) \quad | < \phi | V_{m-nt} | \Phi_m > |^2$$



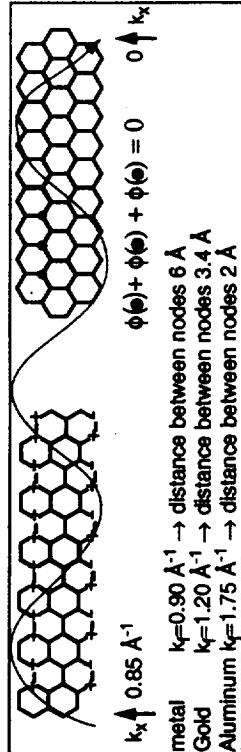
ARMCHAIR	ZIGZAG
$E=0$ at $k_x = 2\pi/3a_0 = 0.85 \text{ \AA}^{-1}$ Metal with $k_F \text{Fermi} < 0.85 \text{ \AA}^{-1}$ couple weakly; threshold $k_F \text{Fermi}$	$E=0$ at $k_x = 0$ No threshold for $k_F \text{Fermi}$



Nodes on the cylinder - Shape of NT wave function



$$\Psi = e^{imk_x a_0 \phi}$$

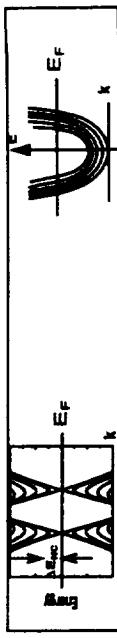


Scattering rate from Metal to Nanotube $\propto \langle \Psi_{\text{nanotube}} | H \text{ coupling} | \Psi_{\text{metal}} \rangle^2$

- Side-contacted: zigzag nanotube are more desirable (curvature)
- Larger metal Fermi wave vector helps.

Conclusions

- dI/dV versus V does not increase in a manner commensurate with the increase in number of subbands.



- The increase in dI/dV with bias is much smaller than the increase in the number of subbands - a consequence of Bragg reflection

- Requirement for axial wave vector conservation:

ZIGZAG cut-off $k_F \text{ Fermi} = 2\pi/3a_0 = 0.85 \text{ \AA}^{-1}$

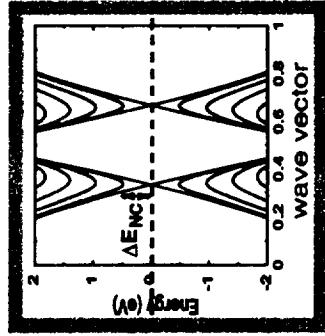
ARMCHAIR cut-off $k_F \text{ Fermi} = 0$

- Our calculations show an increase in transmission with length of contact, as seen in experiments.

- It is desirable for molecular electronics applications to have a small contact area, yet large coupling. In this case, the circumferential dependence of the nanotube wave function dictates:

- Transmission in armchair tubes saturates around unity
 - Transmission in zigzag tubes saturates at two

At what applied voltage are electrons injected into higher subbands?



Bias at which electrons are injected into non crossing subbands is ΔE_{nc}

size	(5,5)	(10,10)	(20,20)	(40,40)
ΔE_{nc} (eV)	1.9	0.98	0.5	0.25

For example, in a (20,20) nanotube electrons are injected into over 20 subbands at an energy of 2.5 eV.

The maximum conductance if the Fermi energy is at 2.5 eV is $\sim 40e^2/h$